

AC Stark effect

Note Title

10/10/2013

"Dipole potential" or "light shift"

Consider the Schrodinger equation for classical light interacting with a quantum atom:

$$i\hbar \frac{\partial}{\partial t} |\psi\rangle = H |\psi\rangle \quad |\psi\rangle = c_g |g\rangle + c_e |e\rangle$$

$$\begin{cases} i\hbar \dot{c}_g = \hbar \frac{\Omega}{2} e^{i\delta t} c_e \\ i\hbar \dot{c}_e = \hbar \frac{\Omega}{2} e^{-i\delta t} c_g \end{cases} \quad \text{(ignoring the phase of the electric field)}$$

using: $\tilde{c}_g = c_g e^{-i\delta t/2} \quad \tilde{c}_e = c_e e^{i\delta t/2}$

the equation for $\dot{\tilde{c}}_e$ becomes:

$$i\hbar \frac{\partial}{\partial t} (e^{-i\delta t/2} \tilde{c}_e) = \hbar \frac{\Omega}{2} e^{-i\delta t/2} \tilde{c}_g$$
$$i\hbar \dot{\tilde{c}}_e e^{-i\delta t/2} + \frac{\delta}{2} \hbar \tilde{c}_e e^{-i\delta t/2} = \hbar \frac{\Omega}{2} e^{-i\delta t/2} \tilde{c}_g$$

similarly:
$$\begin{cases} i\hbar \dot{\tilde{c}}_e = \hbar \frac{\Omega}{2} \tilde{c}_g - \frac{\delta}{2} \hbar \tilde{c}_e \\ i\hbar \dot{\tilde{c}}_g = \hbar \frac{\Omega}{2} \tilde{c}_e + \frac{\delta}{2} \hbar \tilde{c}_g \end{cases}$$

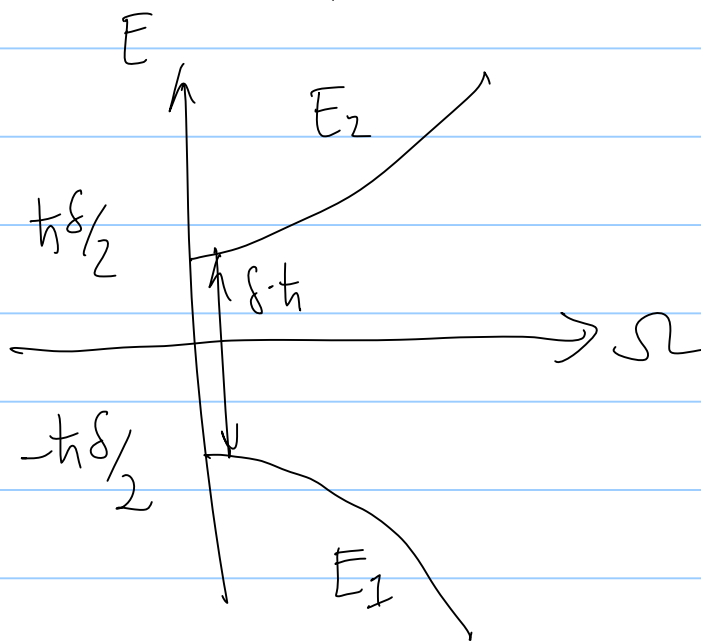
As for any quantum 2-level problem, these equations look like time evolution according to an effective Hamiltonian:

$$\tilde{H} = \frac{\hbar}{2} \begin{pmatrix} \delta/2 & \Omega/2 \\ \Omega/2 & -\delta/2 \end{pmatrix}$$

The energy of our "dressed states" $|1\rangle$ and $|2\rangle$ that diagonalize $\tilde{H}$ are:

$$E_1 = -\frac{\hbar}{2} \sqrt{\delta^2 + \Omega^2} \quad E_2 = \frac{\hbar}{2} \sqrt{\delta^2 + \Omega^2}$$

In the far-detuned limit ($\delta \gg \Omega$), there is no significant excitation, but there are level shifts:

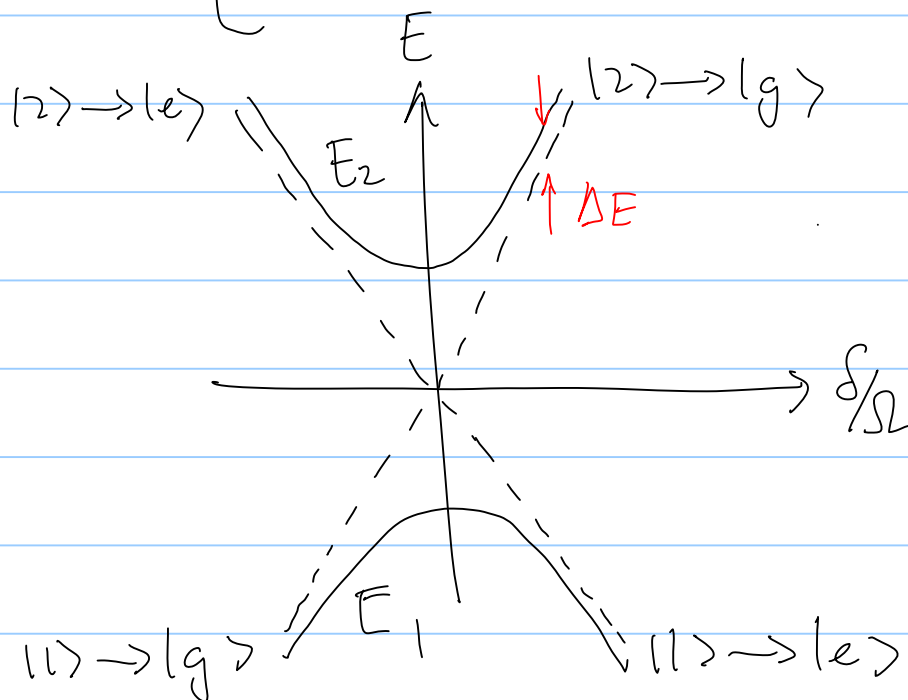


How do we identify the states with $|g\rangle$ and $|e\rangle$?

the eigenvectors of $\tilde{H}$ are:

$$|1\rangle \propto \left[\left(\frac{\delta}{\Omega} - \sqrt{1 + \left(\frac{\delta}{\Omega} \right)^2} \right) |g\rangle + |e\rangle \right]$$

$$|2\rangle \propto \left[\left(\frac{\delta}{\Omega} + \sqrt{1 + \left(\frac{\delta}{\Omega} \right)^2} \right) |g\rangle + |e\rangle \right]$$



Let's now find the energy shift ΔE from the interaction with the light for the ground state $|2\rangle$ for $\frac{\delta}{\Omega} \geq 0$ and $|1\rangle$ for $\frac{\delta}{\Omega} < 0$.

$$@ \Omega=0, E_2 = \frac{\hbar \delta}{2} \text{ and } E_1 = -\frac{\hbar \delta}{2}$$

$$\Delta E_2 = \frac{\hbar}{2} \sqrt{\delta^2 + \Omega^2} - \frac{\hbar \delta}{2}$$

$$= \frac{\hbar}{2} \delta \left[\sqrt{1 + \frac{\Omega^2}{\delta^2}} - 1 \right]$$

$$\approx \frac{\hbar}{2} \delta \left[1 + \frac{1}{2} \frac{\Omega^2}{\delta^2} - 1 \right] \text{ for } |\Omega| \ll |\delta| \text{ or } \left| \frac{\Omega}{\delta} \right| \ll 1$$

$$= \frac{\hbar \Omega^2}{4\delta}$$

Similarly :

$$\Delta E_1 = -\frac{\hbar}{2} \sqrt{\delta^2 + \Omega^2} + \frac{\hbar \delta}{2}$$

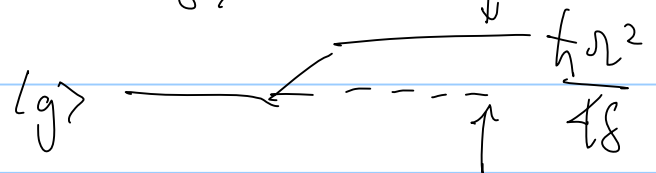
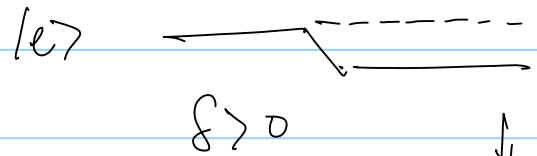
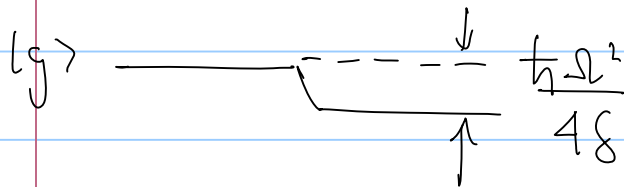
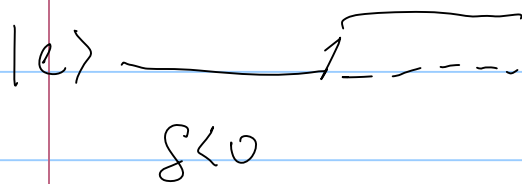
$$= \frac{\hbar}{2} \left(\delta - |\delta| \sqrt{1 + \frac{\Omega^2}{\delta^2}} \right)$$

$$\approx \frac{\hbar}{2} \left(\delta - |\delta| - \frac{1}{2} \frac{\Omega^2}{\delta^2} |\delta| \right) \text{ for } \left| \frac{\Omega}{\delta} \right| \ll 1$$

$$= \frac{-\hbar \Omega^2}{4|\delta|} = \frac{\hbar \Omega^2}{4\delta}$$

(since $\delta < 0$ when $|1\rangle$ is the ground state)

Here's what the light does:



The same analysis for the excited state shows that it shifts by the same amount but in the opposite direction.